

## Sine Gordon model Renormalization

$$L[\psi] = \frac{1}{2v} (\partial_t \psi)^2 - \frac{v}{2} (\partial_x \psi)^2 + \mathcal{G} \cos \beta \psi$$

integrations by part

$$S[\psi] = \int \int dt dx \left[ \underbrace{\frac{v}{2} \psi \partial_x^2 \psi - \frac{1}{2v} \psi (\partial_t \psi)^2}_{S_0[\psi]} + \underbrace{\mathcal{G} \cos \beta \psi}_{S_I[\psi]} \right]$$

split into fast and slow modes:

$$\psi(x) = \psi^s(x) + \delta\psi(x)$$

$$(\psi^s(x) + \delta\psi(x)) (\partial_x^2 \psi^s(x) + \partial_x^2 \delta\psi(x))$$

$$\int dx \delta\psi \nabla_x^2 \psi^s = \int dx \psi^s \nabla_x^2 \delta\psi \quad \xrightarrow{\text{Fourier transform}}$$

$$= \int dx \int_{\text{bulk}} \frac{d\varphi}{(2\pi)^2} \int_{\text{shell}} \frac{d\varphi'}{(2\pi)^2} \psi(\varphi) \left( -\varphi'^2 - \frac{\omega^2}{v^2} \right) \psi(\varphi')$$

$$e^{i(\varphi+\varphi')x}$$

$$= \int dx \int_{\text{bulk}} \frac{d\varphi}{(2\pi)^2} \psi(\varphi) \left( -\varphi^2 - \frac{\omega^2}{v^2} \right) \vartheta(|\varphi| - \frac{1}{s})$$

Definition:

$$\frac{1}{2} \nabla_x^2 G(x, x') = \delta(x - x')$$

$$\frac{1}{2} \nabla_x^2 \int \frac{d^2 q}{(2\pi)^2} G_0(q) e^{i q(x-x')} = \int \frac{d^2 q}{(2\pi)^2} e^{i q(x-x')}$$
$$= \int \frac{d^2 q}{(2\pi)^2} G_0(q) \cdot \frac{1}{2} \left( -q^2 - \frac{\omega^2}{v^2} \right) e^{i q(x-x')}$$

$$G_0(q) = - \frac{2v^2}{q^2 v^2 + \omega^2}$$

Dyson eq:

$$G_T(x, x') = G_0(x, x') - \int \overline{\Sigma}(x, x'') \left. \begin{array}{l} \text{reverse} \\ \downarrow \end{array} \right\} G(x'', x')$$

$$\frac{1}{2} \nabla_x^2 G(x, x') - \int d^2 x'' \overline{\Sigma}(x, x'') G(x'', x')$$
$$= \frac{1}{(2\pi)^2} \delta(x - x')$$

$$\Rightarrow \int \frac{d^2 q}{(2\pi)^2} \frac{d^2 q'}{(2\pi)^2} G(q, q') \cdot \frac{1}{2} \left( -q^2 - \frac{\omega^2}{v^2} \right) e^{i q x + i q' x'} \left. \begin{array}{l} \downarrow \\ \end{array} \right\}$$

$$- \int d^4x'' \int \frac{d^2\tilde{q}}{(2\pi)^2} \frac{d^2\tilde{k}}{(2\pi)^2} \frac{d^2\tilde{k}'}{(2\pi)^2} \frac{d^2\tilde{q}'}{(2\pi)^2} \Sigma(\tilde{q}, \tilde{k}) G(\tilde{k}', \tilde{q}')$$

$$e^{i\tilde{q}x + i(\tilde{k} + \tilde{k}')x'' + i\tilde{q}'x'} = \int \frac{d^2\tilde{q}''}{(2\pi)^2} e^{i\tilde{q}(x-x'')} \int \frac{d^2\tilde{q}'}{(2\pi)^2} e^{i\tilde{q}'(x''-x')}$$

↓  
integrate out

$$\int \frac{d^2\tilde{q}}{(2\pi)^2} \frac{d^2\tilde{q}'}{(2\pi)^2} \left[ G(\tilde{q}, \tilde{q}') G_0^{-1}(\tilde{q}) - \int \frac{d^2\tilde{q}''}{(2\pi)^2} \Sigma(\tilde{q}, \tilde{q}'') \right]$$

$$G(-\tilde{q}'', \tilde{q}') e^{i\tilde{q}x + i\tilde{q}'x'} = \int \frac{d^2\tilde{q}''}{(2\pi)^2} \frac{d^2\tilde{q}'}{(2\pi)^2} \delta(\tilde{q} + \tilde{q}')$$

$$e^{i\tilde{q}x + i\tilde{q}'x'}$$

$$\Rightarrow G(\tilde{q}, \tilde{q}') = G_0(\tilde{q}) \delta(\tilde{q} + \tilde{q}') + G_0(\tilde{q})$$

$$\int \frac{d^2\tilde{q}''}{(2\pi)^2} \Sigma(\tilde{q}, \tilde{q}'') G(-\tilde{q}'', \tilde{q}')$$

- Perturbative theory in powers of  $g$

$$G^{(0)}(\tilde{q}, \tilde{q}') = G_0(\tilde{q}) \delta(\tilde{q} + \tilde{q}') \quad (\text{zero-th order})$$

up to first order

$$G^{(1)}(\tilde{q}, \tilde{q}') = G_0(\tilde{q}) \delta(\tilde{q} + \tilde{q}') - G_0(\tilde{q}) \int \frac{d^2\tilde{q}''}{(2\pi)^2} b^S(\tilde{q}, \tilde{q}'', \tilde{q}') G(-\tilde{q}'', \tilde{q}')$$



$$G^{(1)}(q, q') = G_0(q) \delta(q+q') - \frac{1}{(2\pi)^2} G_0(q) b^S(q, q') G_0(-q')$$

second order

$$G^{(2)}(q, q') = G_0(q) \delta(q+q') - G_0(q) \int \frac{d^2 q''}{(2\pi)^2} b^S(q, q'') \\ \left[ G_0(-q'') \delta(-q''+q') - \frac{1}{(2\pi)^2} G_0(q'') b^S(-q'', q') G_0(-q') \right]$$

---


$$\Rightarrow G^{(2)}(q, q') = G_0(q) \delta(q+q') - \frac{1}{(2\pi)^2} G_0(q) b^S(q, q')$$

$$G_0(-q') + \frac{1}{(2\pi)^4} G_0(q) \int \frac{d^2 q''}{(2\pi)^2} b^S(q, q'') G_0(-q')$$

$$b^S(-q'', q') G_0(-q')$$


---

with

$$\delta S[\psi^S, \psi] = S_0[\psi] + \int dx a^S(x) \delta \psi(x) + \\ \int dx' dx \delta \psi(x) b^S(x, x') \delta \psi(x')$$

$$\rightarrow \int dx dx' \delta \psi(x) \left[ \delta(x-x') \frac{1}{2} \nabla_x^2 + b^S(x, x') \right] \delta \psi(x') \\ + \int dx a^S(x) \delta \psi(x)$$

$$- \int d\vec{x}'' \int \frac{d\vec{q}}{(2\pi)^2} \frac{d\vec{k}}{(2\pi)^2} \frac{d\vec{k}'}{(2\pi)^2} \frac{d\vec{q}'}{(2\pi)^2} \Sigma(\vec{q}, \vec{k}) G(\vec{k}', \vec{q}')$$

$$e^{i\vec{q}\cdot\vec{x} + i(\vec{k} + \vec{k}')\cdot\vec{x}'' + i\vec{q}'\cdot\vec{x}'} = \int \frac{d\vec{q}''}{(2\pi)^4} e^{i\vec{q}(\vec{x} - \vec{x}'')}$$

↓  
integrate out

$$\int \frac{d\vec{q}}{(2\pi)^2} \frac{d\vec{q}'}{(2\pi)^2} \left[ G(\vec{q}, \vec{q}') G_0^{-1}(\vec{q}) - \int \frac{d\vec{q}''}{(2\pi)^2} \Sigma(\vec{q}, \vec{q}'') \right]$$

$$G(-\vec{q}'', \vec{q}') \Big] e^{i\vec{q}\cdot\vec{x} + i\vec{q}'\cdot\vec{x}'} = \int \frac{d\vec{q}''}{(2\pi)^2} \frac{d\vec{q}'}{(2\pi)^2} \delta(\vec{q} + \vec{q}')$$

$$e^{i\vec{q}\cdot\vec{x} + i\vec{q}'\cdot\vec{x}'}$$

$$\Rightarrow G(\vec{q}, \vec{q}') = G_0(\vec{q}) \delta(\vec{q} + \vec{q}') + G_0(\vec{q})$$

$$\int \frac{d\vec{q}''}{(2\pi)^2} \Sigma(\vec{q}, \vec{q}'') G(-\vec{q}'', \vec{q}')$$

- perturbative theory in powers of  $g$

$$G^{(0)}(\vec{q}, \vec{q}') = G_0(\vec{q}) \delta(\vec{q} + \vec{q}') \quad (\text{zero-th order})$$

up to first order

$$G^{(1)}(\vec{q}, \vec{q}') = G_0(\vec{q}) \delta(\vec{q} + \vec{q}') - G_0(\vec{q}) \int \frac{d\vec{q}''}{(2\pi)^2} b^S(\vec{q}, \vec{q}'') G(-\vec{q}'', \vec{q}')$$

$$G^{(1)}(q, q') = G_0(q) \delta(q+q') - \frac{1}{(2\pi)^2} G_0(q) b^S(q, q') G_0(-q')$$

second order

$$G^{(2)}(q, q') = G_0(q) \delta(q+q') - G_0(q) \int \frac{d^2 q''}{(2\pi)^2} b^S(q, q'') \\ \left[ G_0(-q'') \delta(-q''+q') - \frac{1}{(2\pi)^2} G_0(q'') b^S(-q'', q') G_0(-q') \right]$$

---


$$\Rightarrow G^{(2)}(q, q') = G_0(q) \delta(q+q') - \frac{1}{(2\pi)^2} G_0(q) b^S(q, q') G_0(-q')$$

$$G_0(-q') + \frac{1}{(2\pi)^2} G_0(q) \int \frac{d^2 q''}{(2\pi)^2} b^S(q, q'') G_0(-q')$$

$$b^S(-q'', q') G_0(-q')$$


---

with

$$\delta S[\psi^S, \psi] = S_0[\psi] + \int dx a^S(x) \psi(x) + \\ \int dx dx' \psi(x) b^S(x, x') \psi(x')$$

$$\rightarrow \int dx dx' \psi(x) \left[ \delta(x-x') \cdot \frac{1}{2} \nabla_x^2 + b^S(x, x') \right] \psi(x') \\ + \int dx a^S(x) \psi(x)$$



$$\rightarrow \delta S[\psi^S, \delta\psi] = \int dx dx' \delta\psi(x) G^{-1}(x, x') \delta\psi(x') \\ + \int dx a^S(x) \psi(x)$$

$$\delta\psi(x) = \tilde{\psi}(x) + r(x)$$

$$\delta S[\psi^S, \delta\psi] = \int dx dx' \left\{ \tilde{\psi}(x) G^{-1}(x, x') \tilde{\psi}(x') \right. \\ \left. + \tilde{\psi}(x) \left[ 2G^{-1}(x, x') r(x') + a^S(x') \delta(x-x') \right] \right. \\ \left. + r(x) \left[ G^{-1}(x, x') r(x') + a^S(x') \delta(x-x') \right] \right\}$$

Now if

$$\int dx' G^{-1}(x, x') r(x') = -\frac{1}{2} a^S(x)$$

$$\delta S[\psi^S, \delta\psi] = \int dx dx' \left[ \tilde{\psi}(x) G^{-1}(x, x') \tilde{\psi}(x') \right. \\ \left. - \frac{1}{4} a^S(x) G^{-1}(x, x') a^S(x') \right]$$

$$\Rightarrow \int dx'' \int dx' \underbrace{G(x'', x) G^{-1}(x, x')}_{\delta(x''-x')} \psi(x')$$

$$= r(x'') = -\frac{1}{2} \int dx' G(x, x') a^S(x')$$

$$\psi(x) = \psi^s(x) + \tilde{\psi}(x) + \varphi(x) = \tilde{\psi}(x) + \psi^s(x)$$

$$H_0[\Pi, \Psi] = \int dx \frac{v^2}{2} [\Pi^2 + (\partial_x \Psi)^2] \rightarrow H_0[\Pi^s, \psi^s]$$

$$+ H_0[\tilde{\Pi}, \tilde{\Psi}]$$

unperturbed Green function

$$G_0(x, x') = \langle 0 | \psi(x) \psi(x') | 0 \rangle^{\psi}$$

$$= \langle 0^{\psi} | (\tilde{\psi}(x) + \psi^s(x)) (\tilde{\psi}(x') + \psi^s(x')) | 0^{\psi} \rangle$$

$$\stackrel{\psi}{=} G_0(x, x') = 2 \langle 0^{\psi} | \tilde{\psi}(x) \tilde{\psi}(x') | 0^{\psi} \rangle$$

$$= 2 \langle 0 | \tilde{\psi}(x) \tilde{\psi}(x') | 0 \rangle^{\tilde{\psi}}$$

the effective slow mode residual contribution to the action

$$\mathcal{S}_{\text{eff}}[\psi^s] = \langle 0^{\psi} | \mathcal{S}[\psi^s, \varphi] | 0^{\psi} \rangle$$

$$\mathcal{S}_{\text{eff}}[\psi^s] = \int dx dx' \left[ \frac{1}{2} G_0(x, x') G^{-1}(x', x) \right. \\ \left. - \frac{1}{4} a^s(x) G(x, x') a^s(x') \right]$$



$$\begin{aligned}
8 S_{\text{eff}}[\Psi^0] &= \int d\alpha d\alpha' \left[ \frac{1}{2} G_0(\alpha, \alpha') \left[ G_0(\alpha, \alpha') + b^S(\alpha, \alpha') \right] \right. \\
&\quad \left. - \frac{1}{4} a^S(\alpha) G_0(\alpha, \alpha') a^S(\alpha') \right] \\
&= \langle 0 | S_0[\Psi^0] | 0 \rangle^{\Phi} + \int d\alpha d\alpha' \left[ \frac{1}{2} G_0(\alpha, \alpha') b^S(\alpha, \alpha') \right. \\
&\quad \left. - \frac{1}{4} a^S(\alpha) G_0(\alpha, \alpha') a^S(\alpha') \right] \leftarrow 8 S_{\text{eff}}[\Psi^0]
\end{aligned}$$

Let's compute first term:

$$\begin{aligned}
F_1[\Psi^0] &= \frac{1}{2} \int d\alpha d\alpha' \int_{\text{Sheet}} \frac{d^2 q}{(2\pi)^2} \frac{d^2 q'}{(2\pi)^2} G_0(q) \\
&\quad b^S(q', q'') e^{i q(\alpha - \alpha') + i q' \alpha + i q'' \alpha'} \\
&\quad \downarrow \text{integrate out } \alpha, \alpha' \\
&= \frac{1}{2} \int_{\text{Sheet}} \frac{d^2 q}{(2\pi)^2} G_0(q) \theta' b^S(-q, q)
\end{aligned}$$

---


$$\begin{aligned}
b^S(\alpha, \alpha') &= -\frac{\beta^2}{2} g \cos \beta \varphi \delta(\alpha - \alpha') \\
&= -\frac{\beta^2}{2} \mathcal{L}_I[\Psi^0] \delta(\alpha - \alpha')
\end{aligned}$$

$$b^s(p, z') = \int dx dx' b^s(x, x') e^{i\frac{p}{2}x + i'q'x'} \delta(x - x')$$

$$= -\frac{\beta^2}{2} \int \cos \beta \Phi^s$$

$$= b^s(q + q')$$

$$\bullet \quad F_1[\Phi^s] = \frac{1}{2} b^s(q=0) \int_{\text{sheet}} \frac{dq^2}{(2\pi)^2} G_0(q)$$

$$= -b^s(q=0) \int_{\text{sheet}} \frac{dq^2}{(2\pi)^2} \frac{1}{q^2 + \omega^2/v^2}$$

$$= -\frac{1}{2\pi} b^s(q=0) \int_{NS}^{\Lambda} \frac{1/q^2}{q^2} dq$$

$$= -\frac{1}{2\pi} b^s(q=0) \log s$$

$$\rightarrow \frac{\beta^2}{4\pi} \log s \int dx F_1[\Phi^s] = -\frac{g^2 \alpha'}{4\pi} b^2 \int dx \cos \beta \Phi^s$$

$$F_2[\Phi^s] = - \int dx dx' \frac{1}{4} G_0(x, x') a^s(x) a^s(x')$$

=

$$S_R[\Psi] = \int dx \left[ \frac{1}{2\beta^2} \left( 1 + \frac{3\beta^4 g^2 dl}{4\pi \Lambda^3} \right) \Psi \partial_x^2 \Psi + g \Psi^{-2} \left( 1 - \frac{\beta^2 dl}{4\pi} \right) \cos \Psi \right]$$

$$S[\Psi] = \int dx \left[ \frac{1}{2\beta^2} \Psi \nabla_x^2 \Psi + g \cos \Psi \right]$$

$$\Rightarrow \begin{cases} \beta_R^{-2} = \beta^{-2} \left( 1 + \frac{3\beta^4 g^2 dl}{4\pi \Lambda^3} \right) \\ g_R = g \Psi^{-2} \left( 1 - \frac{\beta^2 dl}{4\pi} \right) \end{cases}$$

$$\Rightarrow \frac{d\beta^{-2}}{dl} = \frac{3\beta^2 g^2}{4\pi \Lambda^3} = \frac{-2}{\beta^{-3}} \frac{d\beta}{dl}$$

$$\Rightarrow \boxed{\frac{d\beta^2}{dl} = -\frac{3\beta^6 g^2}{4\pi \Lambda^3}}$$

$$g_R = g (1 + 2dg) \left( 1 - \frac{\beta^2 dl}{4\pi} \right) = g + g \left( 2 - \frac{\beta^2}{4\pi} \right) dl$$

$$\Rightarrow \boxed{\frac{dg}{dl} = 2g \left( 1 - \frac{\beta^2}{8\pi} \right)}$$

$$\text{Let: } K = \frac{\beta^2}{8\pi}, \quad \mu = 4 \sqrt{\frac{3\pi}{\Lambda^3}} g$$



$$\frac{du}{dt} = 2u(1-k)$$

$$\frac{dk}{dt} = -u^2 k^3$$

$$k = 1+v \quad \left\{ \begin{array}{l} \frac{du}{dt} = -2uv \\ \frac{dv}{dt} = -u^2(1+2v) \end{array} \right.$$

$$\Rightarrow 2u \frac{du}{dt} = \frac{du^2}{dt} = -4u^2 v$$

$$2v \frac{dv}{dt} = -2u^2 v - \frac{4u^2 v^2}{\downarrow} = \frac{dv^2}{dt}$$

second order

$$\Rightarrow \frac{d}{dt}(u^2 - 2v^2) = \frac{d}{dt}(u^2 - 2(k-1)^2) = 0$$

