

Classical nonlinear sigma model:

$$S_0 = \frac{1}{2g} \int dx^2 \partial^\mu g^{-1} \partial_\mu g$$

variation: $g \rightarrow g + \delta g$

$$\partial^\mu g g^{-1} \partial_\mu g + \partial^\mu g^{-1} \partial_\mu \delta g$$

$$\cdot \partial^\mu g^{-1} \partial_\mu \delta g = -g^{-1} \partial^\mu g g^{-1} \partial_\mu \delta g = -\partial_\mu (g^{-1} \partial^\mu g \delta g) + \partial_\mu (g^{-1} \partial^\mu g g^{-1}) \delta g$$

$$= \partial_\mu (g^{-1} \partial^\mu g) g^{-1} \delta g + g^{-1} \partial^\mu g \delta g$$

$$\cdot \partial^\mu g g^{-1} \partial_\mu g = -\partial^\mu (g^{-1} \delta g g^{-1}) \partial_\mu g = -\partial^\mu (g^{-1} \delta g g^{-1}) \partial_\mu g + g^{-1} \delta g g^{-1} \partial^\mu \partial_\mu g$$

$$\cdot g^{-1} \partial^\mu g \partial_\mu g^{-1} \delta g = -g^{-1} \partial^\mu g g^{-1} \partial_\mu g g^{-1} \delta g = -g^{-1} \delta g g^{-1} \partial^\mu g g^{-1} \partial_\mu g$$

The motion of simple conservation of currents:

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$$Sg = \frac{1}{2g} \int d^2x \left\{ \text{tr} (g^{-1} \partial_\mu g \partial^\mu (g^{-1} \partial^\nu g)) + \text{tr} (g^{-1} \partial_\mu g \partial^\mu (g^{-1} \partial^\nu g)) \right\}$$

• Lemma: $\partial^\mu (g^{-1} \partial_\mu g) \neq \partial_\mu (g^{-1} \partial^\mu g)$

$$\begin{aligned} & \partial^\mu g^{-1} \partial_\mu g + g^{-1} \partial^\mu \partial_\mu g \\ = & g^{-1} \partial^\mu g \partial_\mu g^{-1} g + g^{-1} \partial^\mu \partial_\mu g \end{aligned}$$

~~$\text{tr}(\partial^\mu g)$~~

$$\begin{aligned} \text{tr}(g^{-1} \partial^\mu g g^{-1} \partial_\mu g) &= \text{tr}(\partial^\mu g g^{-1} \partial_\mu g) \\ = & \text{tr}(\partial^\mu g g^{-1} g \partial_\mu g^{-1} g g^{-1} \partial_\mu g) \end{aligned}$$

conclusion: $\partial^\mu (g^{-1} \partial_\mu g) \neq \partial_\mu (g^{-1} \partial^\mu g)$

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The motion of implies conservation of currents:

$$J_\mu = g^{-1} \partial_\mu g \rightarrow \partial^\mu J_\mu = 0$$

if we introduce holomorphic and anti-holomorphic

current:

$$\partial_{\bar{z}} J_{\bar{z}} + \partial_{\bar{z}} J_{\bar{z}} = 0$$

$$\begin{aligned} & (\partial_1 - i \partial_2)(J_1 - i J_2) + (\partial_1 + i \partial_2)(J_1 + i J_2) \\ = 0 & \Rightarrow \partial_1 J_1 + \partial_2 J_2 = 0 \end{aligned}$$

Chiral current: $\epsilon^{\mu\nu} J_\nu = \tilde{J}_\mu$

$$\begin{aligned} \bullet \quad \cancel{D_\mu J_\mu} &= [D_\mu, D_\nu] = \partial_\mu J_\nu - \partial_\nu J_\mu + [J_\mu, J_\nu] \\ &= 0 \end{aligned}$$

$\Rightarrow J_\mu$ is the pure gauge current

WZW model

$$\Gamma = \frac{-i}{24\pi} \int_{\mathcal{B}} d^3 y \operatorname{Exp} \operatorname{tr} (g^{-1} \partial_\alpha g g^{-1} \partial_\beta g g^{-1} \partial_\gamma g)$$

Feynman diagram operation as:

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• then we consider action

$$S = S_0 + K T$$

variation of WZW term:

$$\frac{-i}{24\pi} \int_{\mathcal{B}} \mathcal{L}^{\alpha\beta\gamma} \text{tr} \left(g^{-1} \partial_\alpha g g^{-1} \partial_\beta g g^{-1} \partial_\gamma g \right)$$

$$= -\frac{i}{24\pi} \int_{\mathcal{B}} \mathcal{L}^{\alpha\beta\gamma} \left[\delta(g^{-1} \partial_\alpha g) \left[-\partial_\beta g^{-1} \partial_\gamma g - g^{-1} \partial_\beta \partial_\gamma g \right] \right]$$

$$= \frac{i}{8\pi} \int_{\mathcal{B}} \mathcal{L}^{\alpha\beta\gamma} \left[-g^{-1} \delta g g^{-1} \partial_\alpha g + g^{-1} \partial_\alpha \delta g \right]$$

$$\partial_\beta \left[g^{-1} \partial_\gamma g \right]$$

$$= \frac{i}{8\pi} \int_{\mathcal{B}} \mathcal{L}^{\alpha\beta\gamma} \left[\partial_\alpha \left[g^{-1} \partial_\alpha \delta g \right] \right] \partial_\beta \left[g^{-1} \partial_\gamma g \right]$$

$$= \frac{i}{8\pi} \int_{\mathcal{B}} \mathcal{L}^{\alpha\beta\gamma} \partial_\alpha \left[g^{-1} \partial_\alpha \delta g \partial_\beta \left[g^{-1} \partial_\gamma g \right] \right]$$

$$= \frac{i}{8\pi} \int_{\mathcal{S}^2} g^{-1} \partial_\alpha \delta g \partial_\beta \left[g^{-1} \partial_\gamma g \right]$$

• Full motion equation is:

$$\partial^\mu (\partial_\mu g^{-1} \partial_\nu g) + \frac{g k}{4\pi} \epsilon^{\mu\nu} \partial^\mu (g^{-1} \partial_\nu g) = 0$$

$$\Rightarrow \left(1 + \frac{g k}{8\pi}\right) \partial_z (g^{-1} \partial_{\bar{z}} g) + \left(1 - \frac{g k}{8\pi}\right) \partial_{\bar{z}} (g^{-1} \partial_z g) = 0$$

= 0

Let: $g = \frac{4\pi}{k}$, the current conservation law

$$\partial_z (g^{-1} \partial_{\bar{z}} g) = 0$$

$$\Rightarrow S_{MNS} = \frac{k}{16\pi} \int d^2x (\partial^\mu g^{-1} \partial_\mu g) + k T$$

• Current algebra

$$J(z) = -k \partial_z g g^{-1}$$

$$\bar{J}(\bar{z}) = k g^{-1} \partial_{\bar{z}} g$$

$$\delta S = -\frac{1}{2\pi} \int d^2x \left[\partial_{\bar{z}} \text{Tr}(\omega(z) J(z)) + \partial_z \text{Tr}(\omega(\bar{z}) \bar{J}(\bar{z})) \right]$$

$$\int d^2x \rightarrow \int \frac{-i}{2} dz d\bar{z}$$

(2)

$$w = \partial_z \psi - [J, \psi] =$$

(6)

where: $J = \sum_a J^a t^a$ $\omega = \sum_a \omega^a t^a$

$$\delta_{\mu, \bar{\omega}} S = \frac{i}{4\pi} \oint dz \text{Tr}(\omega(z) \psi J(z)) - \frac{i}{4\pi} \oint d\bar{z} \text{Tr}(\bar{\omega}(\bar{z}) \bar{J}(\bar{z}))$$



$$\delta_{\mu, \bar{\omega}} S = -\frac{1}{2\pi i} \oint dz \sum_a \omega^a t^a + \frac{1}{2\pi i} \oint d\bar{z} \sum_a \bar{\omega}^a \bar{J}^a$$

ward identity

$$\delta_{\mu, \bar{\omega}} \langle X \rangle = -\frac{1}{2\pi i} \oint dz \sum_a \omega^a \langle J^a X \rangle + \frac{1}{2\pi i} \oint d\bar{z} \sum_a \bar{\omega}^a \langle \bar{J}^a X \rangle$$

we calculate $\delta_{\mu} J$

$$\begin{aligned} \delta_{\mu} J &= -k \delta_{\mu} (\cancel{\partial_z} \partial_z g g^{-1}) \\ &= -k [\partial_z (\delta_{\mu} g) g^{-1} - \partial_z g g^{-1} \delta_{\mu} g g^{-1}] \\ &= -k [(\partial_z \omega g + \omega \partial_z g) g^{-1}] + k \partial_z g g^{-1} \omega \end{aligned}$$

$$a \cdot b \cdot m \sum_{\mathbb{Z}} = m \quad a \cdot b \cdot l \sum_{\mathbb{Z}} = l \quad : \text{arbitrary } m$$

⑨

$$= [w, J] - k \partial_z w$$

⑦

It can be written as:

$$\partial_w J^a = i \sum_{b,c} f_{abc} w^b J^c - k \partial_z w^a$$

⇒ OPE expansion: ?

$$J^a(z) J^b(w) \sim \frac{k \delta^{ab}}{(z-w)^2} + \frac{i}{c} f_{abc} \frac{J^c(w)}{z-w}$$

• Affine Kac-Moody algebra

$$J^a = \sum_{\eta \in \mathbb{Z}} J_{\eta}^a z^{-\eta-1}$$

$$[J_m^a, J_n^a] = \frac{1}{(2\pi i)^2} \left(\oint_{|z|>|w|} dz \oint dw - \oint_{|z|<|w|} dz \oint dw \right)$$

$$z^m w^n J^a(z) J^b(w)$$

$$= \frac{1}{(2\pi i)^2} \oint_0 dw \oint dz z^m z^n \left(\frac{k \delta^{ab}}{(z-w)^2} + \frac{i f_{abc} J^c(w)}{(z-w)} \right)$$

$$= \frac{1}{2\pi i} \oint_0 dw \left(k m \delta^{ab} w^{\eta+m-1} + i f_{abc} J^c(w) w^{m+n} \right)$$

$$= k m \delta^{ab} \delta_{m+n,0} + i f_{abc} J_{m+n}^c \quad (\text{Kac-Moody Algebra})$$

$$\int_{\gamma} \frac{\gamma \wedge \bar{z}}{z} = \dots$$

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⑧ Sugawara construction

We have to show that the energy momentum tensor satisfies the Virasoro algebra

$$T(z)T(w) \sim \frac{c}{2(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{2T'(w)}{z-w}$$

classically:

$$\Rightarrow T_{\mu\nu} = \partial_\mu \frac{\partial L}{\partial (\partial^\nu \phi)} - \partial_\nu L$$

$$T = \frac{1}{2\alpha'} \delta_{ab} \dot{X}^a \dot{X}^b$$

As the first step, we compute the OPE of

$T(z)$ and $J^a(w)$

$$J^a(z) T(w) = \frac{\gamma}{2\pi\alpha'} \oint_{\gamma} \frac{1}{z-w} J^a(z)$$

$$J^b(x) J^b(w)$$



$$\begin{aligned}
 &= \frac{\gamma}{2\pi i} \oint_w \frac{1}{z-w} \left(J^d(z) J^b\left(\frac{w}{x}\right) J^b(w) + J^b(x) T^b\left(\frac{z}{w}\right) J^d(w) \right) \\
 &= \frac{\gamma}{2\pi i} \oint_w \frac{1}{z-w} \left[\frac{k \delta^{ab}}{(z-x)^2} + \frac{i f_{abc} J^c(x)}{z-x} \right] T^b(w) \\
 &+ \frac{\gamma}{2\pi i} \oint_w \frac{J^b(x)}{z-w} \left[\frac{k \delta^{ab}}{(z-w)^2} + \frac{i f_{abd} J^d(w)}{z-w} \right] \\
 &= \frac{\gamma}{2\pi i} \oint_w \frac{1}{z-w} \left[\frac{k \delta^{ab}}{(z-x)^2} + \frac{k \delta^{ab}}{(z-w)^2} \right] T^b(w) \\
 &+ \frac{\gamma}{2\pi i} \oint_w \frac{1}{z-w} \left[i f_{abc} \left[\frac{\delta^{cb}}{(x-w)^2} + \frac{i f_{abd} J^d(w)}{x-w} \right] \frac{1}{z-x} \right. \\
 &+ \left. \frac{\gamma}{2\pi i} \oint_w \frac{1}{z-w} i f_{abc} \left[\frac{\delta^{cb}}{(x-w)^2} + \frac{i f_{abd} J^d(w)}{(x-w)} \right] \frac{1}{z-w} \right] \\
 &= \gamma \left[\frac{2k \delta^{ab} J^b(w)}{(w-x)^2} - \frac{f_{abc} f_{abd} J^d(w)}{(w-x)^2} \right]
 \end{aligned}$$

Dual Coxeter number h^\vee

$$[t^a, t^b] = i f_{abc} t^c$$

$$\Rightarrow [t^b, [t^a, t^c]] = i [t^b, t^c] f_{abc} = i f_{abc} f_{bcd} t^d$$

$$= \frac{1}{2\pi i} \oint \frac{2\pi i (k+h\nu)}{x-w} \left\{ \frac{J^a(x)}{(z-x)^2} + \frac{\partial J^a(x)}{z-x} \right\} - \frac{1}{2\pi i} \oint \frac{2\pi i (k+h\nu)}{(z-x)(x-w)} \left\{ \frac{J^a(x)}{(z-x)^2} + \frac{\partial J^a(x)}{z-x} \right\}$$

(b)

$$\text{tr}(t^d, [t^b, [t^a, t^b]]) = \text{tr}(t^d + t^b + t^a + t^b) - \text{tr}(t^d + t^b + t^b + t^a)$$

$$\Rightarrow T(z) J^a(w) = 2\gamma(k+h\nu) \left(\frac{J^a(w)}{(z-w)^2} + \frac{\partial J^a(w)}{z-w} \right)$$

$$\Rightarrow \gamma = \frac{1}{2(k+h\nu)}$$

• demonstrate $T(z)T(w)$ obey to Virasoro algebra:

$$\begin{aligned} T(z)T(w) &= \frac{1}{4\pi i(k+h\nu)} \oint \frac{dx}{x-w} \left[T(z) J^a(x) J^a(w) \right. \\ &+ \left. T(z) J^a(x) J^a(w) \right] \\ &= \frac{1}{4\pi i(k+h\nu)} \oint \frac{dx}{x-w} \left[\left[\frac{J^a(x)}{(z-x)^2} + \frac{\partial J^a(x)}{z-x} \right] J^a(w) \right. \\ &+ \left. J^a(x) \left[\frac{J^a(w)}{(z-w)^2} + \frac{\partial J^a(w)}{z-w} \right] \right] \end{aligned}$$

$$\begin{aligned}
&= \frac{1}{4\pi i(k+h^v)} \oint \frac{dx}{x-w} \left\{ \frac{R \operatorname{dmg}}{(z-x)^2(x-w)^2} - \frac{2R \operatorname{dmg}}{(z-x)(x-w)^3} \right. \\
&\quad \left. + \frac{R \operatorname{dmg}}{(x-w)^2(z-w)^2} + \frac{2R \operatorname{dmg}}{(z-x)(x-w)^3} \right\} \\
&+ \frac{1}{4\pi i(k+h^v)} \oint \frac{dx}{x-w} \left(\frac{J^a J^a(w)}{(z-x)^2} + \frac{2J^a J^a(w)}{z-x} \right. \\
&\quad \left. + \frac{J^a J^a(w)}{(z-w)^2} + \frac{J^a \partial J^a(w)}{z-w} \right) \\
&= \frac{R \operatorname{dmg}}{(2k+h^v)} + \frac{1}{(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w}
\end{aligned}$$