

operator representation

Matsubara Green's function

$$\tau > 0 : \chi(\tau) = - \langle O_2(\tau) O_1(0) \rangle = \frac{1}{Z}$$

$$O_1(\tau) = e^{H\tau} O_1 e^{-H\tau}$$

$$\text{Log } Z = e^{-\beta H}$$

$$= - e^{\beta H} \sum_{n,m} \langle n | O_2(\tau) | m \rangle \langle m | O_1 | n \rangle$$

$$= - e^{\beta H} \sum_{n,m} e^{-\beta E_n} e^{(E_n - E_m)\tau} \langle n | O_2 | m \rangle \langle m | O_1 | n \rangle$$

\swarrow \downarrow \searrow
 $n \neq m$ \downarrow $E_n > E_m$ \downarrow n

$$\chi(i\omega_n) = - e^{\beta H} \int_0^{\beta} dt e^{i\omega_n t + (E_n - E_m)t} \langle n | O_2 | m \rangle$$

$$\langle m | O_1 | n \rangle e^{-\beta E_n}$$

$$= e^{\beta H} \sum_{n,m} \frac{\langle n | O_2 | m \rangle \langle m | O_1 | n \rangle}{i\omega_n + E_n - E_m} (e^{-\beta E_n} - e^{-\beta E_m})$$

$$= e^{\beta \mu} \sum_{n,m} \frac{\langle n | O_2 | m \rangle \langle m | O_1 | n \rangle}{i\omega_n + E_n - E_m}$$

可以与物理可观测量
建立联系

$$D_{ret}(t-t') = -i \Theta(t-t') \langle [\rho_2(t), \rho_1(t')]_T \rangle$$

$$= -i \Theta(t-t') e^{\beta \mu} \sum_{n,m} \langle n | O_2 | m \rangle \langle m | O_1 | n \rangle$$

↓
热平均

$$e^{i(E_n - E_m)(t-t')} (e^{-\beta E_n} \mp e^{-\beta E_m})$$

$$D_{ret}(\omega) = \int_{-\infty}^{+\infty} dt e^{i\omega t} D_{ret}(t-t')$$

$$= e^{\beta \mu} \sum_{n,m} \frac{\langle n | O_2 | m \rangle \langle m | O_1 | n \rangle}{\omega + (E_n - E_m) + i\eta} (e^{-\beta E_n} \mp e^{-\beta E_m})$$

time-order Green — path integral

• 为时+序 Green 云云

因为 Lehman rep 前向多计算 可观测量 Green

只需要: $i\omega_n \rightarrow \omega + i\eta$

$$G(t-t') = -i \Theta(t-t') \langle | O_2(t) O_1(t') | 0 \rangle$$

$$\mp i \Theta(t-t') \langle | O_1(t') O_2(t) | \rangle$$



零温可以用, 有限温就不行了

$$= -i \Theta(t-t') e^{\beta \mu} \sum_{n,m} \langle n | O_2 | m \rangle \langle m | O_1 | n \rangle e^{-\beta E_n} e^{i(E_n - E_m)(t-t')}$$

$$\mp (-i \Theta(t-t')) e^{\beta \mu} \sum_{n,m} \langle m | O_2 | m \rangle \langle m | O_1 | n \rangle e^{-\beta E_m} e^{i(E_n - E_m)(t-t')}$$

$$= e^{\beta \mu} \sum_{n,m} \langle n | O_2 | m \rangle \langle m | O_1 | n \rangle e^{i(E_n - E_m)(t-t')}$$

$$\left(-i \Theta(t-t') e^{-\beta E_n} e^{-\eta(t-t')} \pm -i \Theta(t-t') e^{-\beta E_m} e^{\eta(t-t')} \right)$$



$$G(\omega) = \sum_{n,m} e^{\beta \mu} \langle n | O_2 | m \rangle \langle m | O_1 | n \rangle$$

$$\left\{ \frac{e^{-\beta E_n}}{\omega + E_n - E_m + i\eta} \mp \frac{e^{-\beta E_m}}{\omega + E_n - E_m - i\eta} \right\}$$

$$G(t-t') = -i \Theta(t-t') \langle | O_2(t) O_1(t') | 0 \rangle$$

$$F: -i \Theta(t-t') \langle | O_1(t') O_2(t) | \rangle$$



零温可以用, 有限温就不行了

$$= -i \Theta(t-t') e^{\beta \Omega} \sum_{n,m} \langle n | O_2 | m \rangle \langle m | O_1 | n \rangle e^{-\beta E_n}$$

$$e^{i(E_n - E_m)(t-t')}$$

$$F: (-i \Theta(t-t')) e^{\beta \Omega} \sum_{n,m} \langle m | O_2 | n \rangle \langle n | O_1 | m \rangle e^{-\beta E_m}$$

$$e^{i(E_n - E_m)(t-t')}$$

$$= e^{\beta \Omega} \sum_{n,m} \langle n | O_2 | m \rangle \langle m | O_1 | n \rangle e^{i(E_n - E_m)(t-t')}$$

$$\left(-i \Theta(t-t') e^{-\beta E_n} e^{-i\eta(t-t')} \pm -i \Theta(t-t') e^{-\beta E_m} e^{i\eta(t-t')} \right)$$



$$G(\omega) = \sum_{n,m} e^{\beta \Omega} \langle n | O_2 | m \rangle \langle m | O_1 | n \rangle$$

$$\left\{ \frac{e^{-\beta E_n}}{\omega + E_n - E_m + i\eta} \mp \frac{e^{-\beta E_m}}{\omega + E_n - E_m - i\eta} \right\}$$

• Matsubara

$$y = (\chi\omega_n) = e^{\beta n} \frac{\langle n | O_2 | m \rangle \langle m | O_1 | n \rangle}{\sum_{n,m} \chi\omega_n + E_n - E_m}$$

$$(e^{-\beta E_n} \mp e^{-\beta E_m}) \quad (\text{time order})$$

Retard Green function 的奇点. 都在实轴的下面

time-order 都在实轴下面 (finite temp)

• Zero temperature : $\beta \rightarrow \infty$

$$G(\omega) = \sum \frac{\langle 0 | O_2 | m \rangle \langle m | O_1 | 0 \rangle}{\omega + E_0 - E_m + i\eta} \mp \frac{\langle m | O_2 | 0 \rangle \langle 0 | O_1 | m \rangle}{\omega + (E_0 - E_m) - i\eta}$$

$$G(\omega) \Big|_{\beta \rightarrow \infty} : y(\chi\omega_n) = \left. \begin{array}{l} i\omega_n \rightarrow \omega + i\text{sgn}(\omega)\eta \end{array} \right|$$

Matsubara Green func $\xrightarrow{\text{解析延拓}}$ Green func

$$\gamma = \gamma_0, \gamma_1 = \gamma_0, \gamma_2 = \gamma_0$$

justin: vector:

Spin orbit coupling

上海交大: Tony Leggett

• Spectral function

$$A_+(\omega) = 2\pi \sum_{n,m} e^{\beta E_n} \langle n | O_2 | m \rangle \langle m | O_1 | n \rangle e^{-\beta E_m} \delta(\omega + E_n - E_m)$$

$$\begin{aligned} A_-(\omega) &= 2\pi \sum_{n,m} e^{\beta E_n} \langle n | O_2 | m \rangle \langle m | O_1 | n \rangle e^{-\beta E_m} \delta(\omega + E_n - E_m) \\ &= 2\pi \sum_{n,m} e^{\beta E_n} \langle m | O_2 | n \rangle \langle n | O_1 | m \rangle e^{-\beta E_n} \delta(\omega - (E_n - E_m)) \end{aligned}$$

$$A(\omega) = -2 \operatorname{Im} \operatorname{Den}(\omega) = A_+(\omega) \pm A_-(\omega)$$

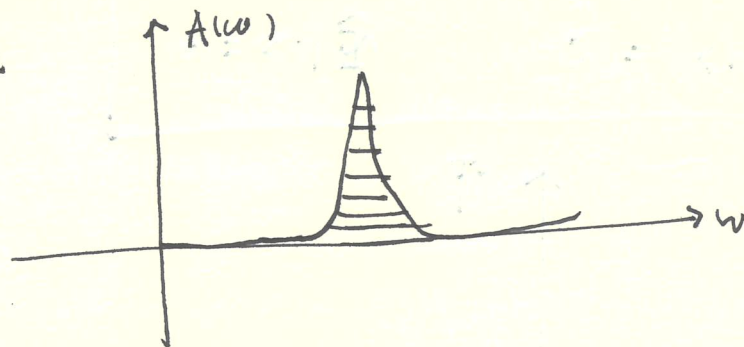
$$A_+(\omega) = e^{\beta \omega} A_-(\omega)$$

if $O_1 = c^\dagger, O_2 = c$

$$\int \frac{d\omega}{2\pi} A(\omega) \equiv \sum_{n,m} e^{\beta E_n} (|\langle n|c|m\rangle|^2 e^{-\beta E_n} + |\langle n|c^\dagger|m\rangle|^2 e^{-\beta E_n}) + e^{-\beta E_n} \langle n|c|m\rangle \langle m|c^\dagger|n\rangle + e^{-\beta E_n} \langle n|c^\dagger|m\rangle$$

$\langle m|c|n\rangle$

$= 1$



谱函数为常数

Bosonic 的谱函数没有这种形式

特征函数:

$$A_+(\omega) = - \frac{2 \text{Im} G(\omega)}{e^{-\beta\omega/2} \pm e^{-\beta\omega}} e^{\pm\beta\omega/2}$$

$$A_+(\omega) - A_-(\omega) = \cot \frac{\beta\omega}{2} \text{Im} D_{ret}(\omega)$$

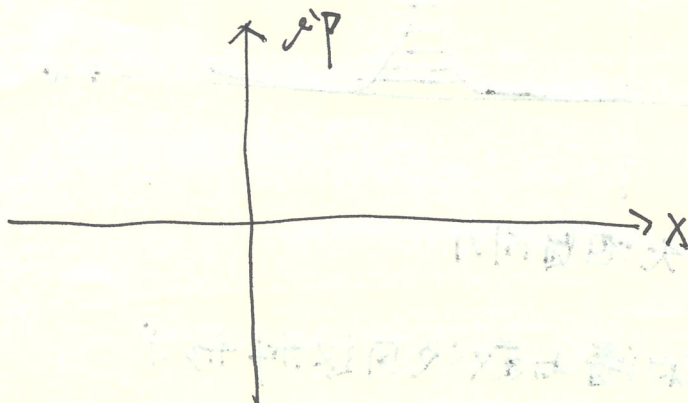
$$\hat{N}^2 = \hat{N}(\hat{N} + 1) = \langle \hat{N} | \hat{N} \rangle$$

Functional field path integral

你不要问我 Grassmann Number 是多少?

Operator: $\{\hat{\Psi}^+, \hat{\Psi}\} = 1 \quad \{\hat{\Psi}, \hat{\Psi}\} = \{\hat{\Psi}^+, \hat{\Psi}^+\} = 0$

$\hat{\Psi}^+ 0\rangle = 1\rangle$	$\hat{\Psi} 1\rangle = 0\rangle$
$\hat{\Psi}^+ 1\rangle = 0$	$\hat{\Psi} 0\rangle = 0$



$$|\bar{\Psi}\rangle = |\alpha\rangle - \hat{\Psi} |1\rangle = e^{-\hat{\Psi} C^+} |0\rangle$$

$$\hat{\Psi} |\bar{\Psi}\rangle = \hat{\Psi} (|\alpha\rangle - \hat{\Psi} |1\rangle)$$

$$\langle \bar{\Psi} | = \langle 0 | - \langle 1 | \hat{\Psi} = \langle 0 | + \hat{\Psi} |1\rangle$$

$$\langle \bar{\Psi} | \bar{\Psi} \rangle = \langle \bar{\Psi} | \bar{\Psi} \rangle$$