

• Resolution identity

$$\int |\psi\rangle\langle\psi| e^{-\bar{\psi}\psi} d\bar{\psi} d\psi = 1$$

Proof:

$$\int |\psi\rangle\langle\psi| e^{-\bar{\psi}\psi} d\bar{\psi} d\psi = 1$$

$$= \int (|0\rangle\langle 0| - |\psi\rangle\langle\psi|) (|0\rangle\langle 0| - |\psi\rangle\langle\psi|) (1 - \bar{\psi}\psi) d\bar{\psi} d\psi$$

$$= \int \bar{\psi}\psi (|0\rangle\langle 0| - |\psi\rangle\langle\psi|) + (|0\rangle\langle 0| - |\psi\rangle\langle\psi|) \bar{\psi}\psi d\bar{\psi} d\psi$$

$$= |0\rangle\langle 0| + |\psi\rangle\langle\psi|$$

• trace identity:

$$\text{tr}(\hat{O}) = \int \langle -\bar{\psi} | \hat{O} | \psi \rangle e^{-\bar{\psi}\psi} d\bar{\psi} d\psi$$

$$= \int \langle 0 | \hat{O} | 0 \rangle - \langle 0 | \hat{O} | \psi \rangle + \langle \psi | \hat{O} | 0 \rangle - \langle \psi | \hat{O} | \psi \rangle d\bar{\psi} d\psi$$

$$= \langle 0 | \hat{O} | 0 \rangle + \langle \psi | \hat{O} | \psi \rangle$$

Fermion many body:

single particle:

$$Z = \text{tr}(e^{-\beta H}) = \int \langle -\Psi_N | e^{-H\beta\Delta\tau} | \Psi_N \rangle$$

$$e^{-\bar{\Psi}_N \Psi_N} d\bar{\Psi}_N d\Psi_N \dots e^{-\bar{\Psi}_{N-1} \Psi_{N-1}}$$

$$= \int \langle -\Psi_N | e^{-H\beta\Delta\tau} | \Psi_{N-1} \rangle \langle \Psi_{N-1} | e^{-H\beta\Delta\tau} | \Psi_{N-2} \rangle$$

$$\dots \langle \Psi_0 | e^{-H\beta\Delta\tau} | \Psi_N \rangle$$

$$\langle \Psi_n | 1 - \beta H \cdot \Delta\tau | \Psi_{n-1} \rangle = e^{-\bar{\Psi}_n \Psi_{n-1}}$$

$$(1 - H(\bar{\Psi}_i \Psi_j))$$

$$= \int e^{-\bar{\Psi}_n (\Psi_n - \Psi_{n-1}) - H(\bar{\Psi}_i - \Psi_j)}$$

$$d\bar{\Psi}_f d\Psi_i$$

$$= \int D[\Psi] e^{\bar{\Psi} S}$$

$$S = \bar{\Psi} \partial_\tau \Psi - H = \bar{\Psi} (\partial_\tau + H) \Psi$$

$$\Psi(\alpha) = \Psi(\beta) \quad ; \quad \Psi(\beta) = -\Psi(\alpha)$$

• Fourier transformation:

$$\Psi(\tau) = \frac{1}{\sqrt{\beta}} \sum_n e^{-i\omega_n \tau} \Psi(i\omega_n)$$

$$\omega_n = \frac{(2n+1)\pi}{\beta}$$

定論中

$$\Psi(\tau) \xrightarrow{\quad} \Psi(\alpha(\tau))$$

$$Z = \int D[\bar{\Psi}(x, \tau), \Psi(x, \tau)] e^{-S[\bar{\Psi}, \Psi]}$$

$$S[\bar{\Psi}, \Psi] = \int_0^\beta d\tau \left(\bar{\Psi}(x, \tau) \partial_\tau \Psi(x, \tau) + H(x, \tau) - \mu \bar{\Psi} \Psi \right)$$

• Bosonic system:

$$\hat{a} |a\rangle = a |a\rangle$$

$$|a\rangle = e^{-\frac{|a|^2}{2}} \sum_n \frac{1}{n!} a^n |0\rangle$$

↑ expand

金屬的蒸發效應

4.

軟物質不是 Quantum

$$\int \frac{dada^*}{\pi} |a\rangle \langle a| = 1$$

$$\langle a|a'\rangle = e^{-\frac{1}{2}(|a|^2 + |a'|^2) + a^* a}$$

trace/identity:

$$\text{tr}(\hat{O}) = \int \langle a|\hat{O}|a\rangle \frac{da da^*}{\pi}$$

Partition function for bosonic system

$$\mathcal{Z} = \text{tr}(e^{-\beta H}) = \int \langle a_N| e^{-\beta H} |a_1\rangle$$

$$\dots \dots \langle a_1| e^{-\beta H} |a_N\rangle$$

=

• Source field:

$$\int d\bar{z}_1 dz_1 \dots d\bar{z}_n dz_n e^{-\bar{z}_i A_{ij} z_j + \bar{\omega}_i z_i + \bar{z}_j \omega_j}$$

$$= \pi^n (\det A^{-1}) e^{\bar{\omega}_i A_{ij} \omega_j}$$

$$\Rightarrow \langle \bar{z}_{i_1}^* z_{j_1} \rangle = (A^{-1})_{j_1 i_1}$$

Wick theorem:

complex Number

$$\langle \bar{z}_{i_1}^* \dots \bar{z}_{i_n}^* z_{j_1} \dots z_{j_n} \rangle$$

$$= \sum_p (A^{-1})_{j_1 p_1} (A^{-1})_{j_2 p_2} \dots (A^{-1})_{j_n p_n}$$

$$\int e^{-\bar{\Phi}_i M_{ij} \Phi_j} \dots d\bar{\Phi}_1 d\Phi_1 \dots d\bar{\Phi}_n d\Phi_n$$

$$= \det(M)$$

• Supersymmetry:

$$(\det M)^{-1} \det M = 1$$

• two particle Hilbert space:

$$\begin{pmatrix} 0 \\ 1 \end{pmatrix}_1 \quad \begin{pmatrix} 0 \\ 1 \end{pmatrix}_2$$

$$\downarrow \quad \quad \quad \downarrow$$

左边的 i 乘上右边的 i 等于 -1 吗?

左边有 i, j, k ; 右边 i, j, k

• $\int e^{-\bar{\Psi} M \Psi + \bar{v}_i v_i + v_i \Psi_i} d\bar{\Psi}_1 d\Psi_1 \dots d\bar{\Psi}_n d\Psi_n = \det(M) e^{\bar{v}_i M_{ij} v_j}$

• $\langle \bar{\Psi}_j, \dots, \bar{\Psi}_n, \Psi_n, \dots, \Psi_j \rangle = \sum_P \text{sgn}(P) A_{j, \bar{1}P_1}^{-1} \dots A_{n, \bar{1}P_n}^{-1}$

\swarrow fermion exchange \searrow Grassmann Number
 Direction Number
 muscle

格林函数 / 重武器主要是技术上